

## ARTICLE

# Vibrational Spectra and Density Functional Theory Calculations of Metallotriphenylcorroles

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The infrared absorption and Raman scattering spectra were measured for the metallotriphenylcorroles (MTPCs, M=Cu, Co, Ni, Mn). The ground-state structures and vibrational spectra of MTPCs have been calculated with the density functional theory. The observed Raman and IR bands have been assigned based on the calculation results. Due to the symmetry lowering, the vibrational spectra of MTPCs are much more complex than metalloporphyrins, and several skeletal modes are found strongly coupled to the phenyl vibrations. The relationship between the Raman/IR frequencies and the structures of TPC ring is investigated. It is found that the vibrations involving the  $C_{\alpha}^I C_{\alpha}^I$  stretch and  $C_{\alpha} C_m$  stretch are sensitive to the size of corrole core. In particular, the frequency of  $\nu_5$ , which is assigned to  $C_{\alpha}^I C_{\alpha}^I$  stretch in coupling with the  $C_{\alpha} C_m$  symmetric stretch, increases linearly with the decrease of the corrole core-sizes and may be used as a mark band to evaluate the structural change of the metallocorroles.

**Key words:** Metallocorroles, Raman spectrum, Molecular vibration, Density functional theory

## I. INTRODUCTION

The corroles as a member of the porphyrinoid family are very important multifunctional compounds. Owing to their unique properties in chemistry and photochemistry, the corroles have great application prospects in a wide range of areas from catalyst, electrochemistry, bio-sensors, nonlinear optics, to photodynamic therapy of tumor [1–3]. They are also used as the structural model of corrin, the core component of vitamin B12. In comparison with normal porphyrins and other related macrocycles, an important feature of the corrole ligands is their capability of stabilizing high oxidation states for coordinated metal ions, due to bearing three ammoniac nitrogen atoms [4–7]. As a result of their unique coordination chemistry, the metallocorroles show distinct catalytic properties for the oxidation or hydroxylation of hydrocarbons [8–11]. For these reasons, the studies of corroles have now become one of the most active branches of porphyrin chemistry. Numerous transition-metal and main-group metal complexes of the corroles have been synthesized and spectroscopically characterized [4–18]. It has been found that the periphery substituents have a significant influence on the electrochemical properties and the electronic absorption spectra of metallocorroles (MCors) [14–16]. For instance, the Soret bands of copper corroles with

different triaryl substituents have been found shift as large as 100 nm [16]. It is clear that periphery modification of the metallocorroles can be an effective route to manipulate their electronic and optical properties.

Raman and IR spectra play an important role in the structural and dynamic studies of porphyrin-related compounds. Theoretically, the density functional theory (DFT) calculations have been widely used to study the vibrational spectra of porphyrin-related compounds in recent years, proven useful in revealing information about the structures and vibrational states of these compounds [19–25]. Although vibrational spectra [17, 18] and DFT calculations [26–30] of several metallocorroles have been reported, descriptions of vibrational modes of metallocorroles, especially the meso-substituted metallocorroles, are still not fully understood.

In this work, we report IR and Raman spectroscopic properties of series metal complexes of meso-triphenylcorrole (MTPC, M=Cu, Ni, Co, and Mn) together with DFT calculations to obtain more detailed knowledge about the vibrational states of these complexes.

## II. EXPERIMENTAL AND COMPUTATIONAL METHODS

The synthesis of free-base meso-triphenylcorrole ( $H_3$ TPC) was adapted from the method by Koszarna and Gryko [31]. 200 mL distilled water and 20 mL methanol was mixed in a flask, then 10 mmol freshly

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distill pyrrole and 5 mmol benzoic aldehyde and 4.25 mL hydrochloric acid (36%) were added into the solvent. After stirring for 3 h at room temperature, the reactant was extracted with chloroform. The organic layer was washed twice with distilled water, dried with anhydrous  $\text{Na}_2\text{SO}_4$ , and then filtrated with diatomite. The obtained solution was diluted to 300 mL with chloroform in a flask to react with 1.23 g (5 mmol) chloranil for 1 h at room temperature. The resultant was chromatographic separated on a silica gel column washed with dichloromethane first and then with 1:1 mixed dichloromethane/*n*-hexane flow phase. The green component was collected. After evaporation of the solvent, the black-green target product was obtained.

The metal complexes, MTPC (M=Cu, Ni, Co, Mn), were synthesized according to the method in Ref.[32], *i.e.*, by treating the metal-acetates and free base  $\text{H}_3\text{TPC}$  in hot chloroform/methanol mixture (7:3 in volume). Taking nickel complex (NiTPC) as an instance, the synthesis details are as follows. Reagent-grade chloroform/methanol mixture (10 mL) was brought gently to a reflux on a hot plate and magnetically stirred to 60 °C. The solid  $\text{H}_3\text{TPC}$  (10 mg) was added into the solvent and allowed to dissolve completely for a few minutes. Then excess nickel acetate ( $\text{Ni}(\text{CH}_3\text{COO})_2 \cdot 4\text{H}_2\text{O}$ ) was added into the solution. Reaction proceeded for about 1 h to allow complete conversion from the free base to the nickel complex. After the reaction mixture cooled, equal volume of distilled water was added into the reaction mixture to precipitate the NiTPC species. The product was separated with a centrifuge and washed with water to remove unreacted nickel acetate. The centrifuging and washing procedures repeated several times until the up-layer solution is clear. The obtained solid NiTPC sample was then silica-gel dried in a desiccator. Similar procedures were carried out to prepare the copper, cobalt, and manganese complexes.

IR spectra of MTPC were measured on a Nicolet MAGNA-IR750 infrared spectrometer as KBr pellets. Raman spectra of MTPC solid powders were measured on a Labram-010 micro-Raman spectrometer. The 488.0 nm line of an  $\text{Ar}^+$  laser (Spectra-Physics 163-C12) was used as the excitation source with the power of 7 mW on samples. The time constant of the air-cooled CCD detector was 15 s, and 4 scans were accumulated for each Raman spectrum.

DFT calculations for the structural optimizations of MTPC were carried out by using the Becke's three-parameter hybrid functional (referred as B3LYP) [33]. Standard 6-31G(d) basis sets were used for C, N, and H, while Ahlrichs' TZVP basis sets for Cu, Ni, Co, and Mn [34]. Frequency calculations at the optimized geometries were conducted to confirm the structure as an energy minimum (with all the frequencies to be positive values) as well as to obtain the theoretical IR and Raman spectra. Assignment of individual vibrational frequency has been carried out by inspecting the calcu-

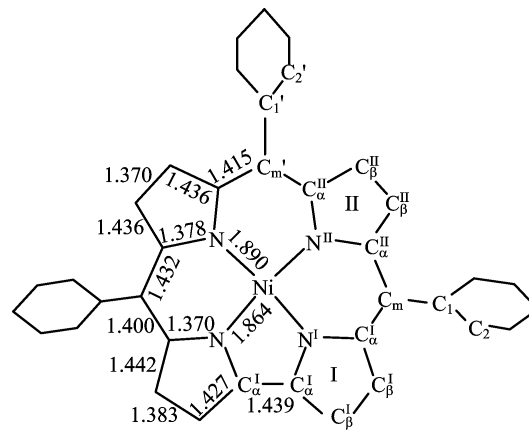


FIG. 1 Structural sketch of NiTPC with atomic labels. The four pyrrole rings are divided into two types (I and II).

lated Cartesian displacements of the corresponding normal mode. As DFT-calculations are known to systematically overestimate harmonic frequencies, the resultant theoretical frequencies are scaled with a single factor of 0.97. The same factor has been used in the DFT calculations of metalloporphyrins and other tetrapyrrole compounds [22–25]. All calculations were performed with the Gaussian 09 [35] program suite on the Lenovo 1800 supercomputer of USTC supercomputation center.

### III. RESULTS AND DISCUSSION

#### A. Geometries and general consideration of the molecular vibrations of MTPCs

Geometry optimizations have been carried out for MTPCs (M=Cu, Ni, Co, and Mn) with various spin-multiplicities. Computational results indicate that the ground state spin-multiplicity is doublet for NiTPC, triplet for CuTPC and CoTPC, and quintet for MnTPC. The ground state structures of all the four complexes are  $C_2$  symmetry with the corrole skeleton to be nearly planar. The maximum derivation of the skeletal atoms from the corrole mean-plane is 0.064 Å for CuTPC, 0.037 Å for NiTPC, 0.032 Å for CoTPC, and 0.043 Å for MnTPC. The M–N bond lengths, which measure the size of the coordination core, are 1.925/1.928 Å for CuTPC, 1.864/1.890 Å for NiTPC, 1.853/1.877 Å for CoTPC, 1.905/1.913 Å for MnTPC.

Figure 1 displays the structure sketch and the atomic labels of NiTPC with some key structural parameters, in which the four pyrrole rings are divided into two sets, denoted as Pyr-I and Pyr-II respectively, with taking the  $C_2$  symmetry into account. The atoms of corrole skeleton are thereby denoted as  $C_\alpha^I$ ,  $C_\beta^I$ ,  $C_\alpha^{II}$ ,  $C_\beta^{II}$ ,  $N^I$ ,  $N^{II}$ ,  $C_m$  (at the two symmetric methines), and  $C'_m$  (at the isolated methine). The phenyl groups of MTPCs are inclined with respect of the corrole mean-plane. For NiTPC, the dihedral angles between the

mean planes of phenyl groups and the corrole skeleton are  $49.6^\circ$  (for  $C_m$ -phenyl groups) and  $64.5^\circ$  (for  $C'_m$ -phenyl). Since the calculated ground state structures and the IR and Raman spectra of the four complexes show similar trends, our discussions in the following sections will be concentrated on NiTPC.

Customarily, a simplified planar molecular model, in which all peripheral substituents are represented by point masses, is often used in the vibrational analyses of the porphyrins and other tetrapyrrole macrocycles for convenience [19–25, 36, 37]. The molecular vibrations of a tetrapyrrole compound can thus be divided into the macrocyclic skeletal modes and the substituent internal modes despite the fact that the motions of the skeletal atoms and the substituent atoms may correlate with each other. The vibrations of macrocyclic skeleton can be further divided into the in-plane (ip) and the out-of-plane (oop) types. With this strategy, the skeleton of a metalcorrole (MCor) can be assumed having a planar  $C_{2v}$  structure and its skeletal modes can be classified as:

$$\Gamma_{\text{MCor-ip}} = 34A_1 + 33B_2$$

$$\Gamma_{\text{MCor-oop}} = 15A_2 + 17B_1$$

For MTPC which has a reduced  $C_2$  symmetry, the skeletal modes should be classified according to the irreducible representations of the  $C_2$  point group:

$$\Gamma_{\text{MTPC-ip}} = 34A + 33B$$

$$\Gamma_{\text{MTPC-oop}} = 15A + 17B$$

The phenyl internal modes, according to the atomic motions in or out of the phenyl planes, can be classified as:

$$\Gamma_{\text{Ph-ip}} = 30A + 30B$$

$$\Gamma_{\text{Ph-oop}} = 14A + 16B$$

In the following sections, the IR/Raman bands of MTPC will be assigned to local coordinates in which the phasing of adjacent bond stretches within the pyrrole rings and at the meso-carbon bridges are taken into account [21, 25, 36, 37]. The local coordinates of MTPC in this work follow those used in metalloporphyrins as defined by Li *et al.* [36, 37]. For in-plane vibrations, the stretches of  $C_m C_1 / C'_m C'_1$ ,  $C'_\beta H / C''_\beta H$ , and  $C^I_\beta C^I_\beta / C^II_\beta C^II_\beta$  bonds and the in-plane bending of  $C^I_\beta H / C^II_\beta H$  are defined as isolated local coordinates. The stretch of  $C^I_\alpha C^I_\alpha$  bond that connects the adjacent Pyr-I rings is also defined as an isolated local coordinate. The stretches of adjacent  $C_\alpha C_m$  bonds are defined as symmetric and asymmetric, respectively, depending on in-phase or out-of-phase. As the stretches of  $C^I_\alpha C^I_\beta / C^II_\alpha C^II_\beta$  and  $C^I_\alpha N^I / C^II_\alpha N^II$  bonds in the same pyrrole ring are heavily coupled with each other, the corresponding vibrations are defined as pyrrole quartering, symmetric half-ring, asymmetric half-ring, and pyrrole breathing coordinates. The vibrations involving mainly the bending of the pyrrole bond-angles are defined as pyrrole symmetric and asymmetric deformations. For the out-of-plane vibrations, local coordinates are defined as pyrrole symmetric or asymmetric folding, pyrrole swivel, pyrrole tilting *etc.* The atomic mo-

tions for these local coordinates can be found in Ref. [36, 37]. The vibrations of phenyl groups can be classified into two types, in-plane and out-of-plane modes, respectively, each of which is further divided into symmetric or asymmetric with respect to the reflection of the plane that is perpendicular to the phenyl plane and passes the  $C_m C_1 / C'_m C'_1$  bond. The phenyl modes are thereby denoted as  $\varphi$  (in-plane symmetric),  $\psi$  (in-plane asymmetric),  $\pi$  (out-of-plane symmetric), and  $\sigma$  (out-of-plane asymmetric).

The calculated vibrational frequencies of NiTPC and their local coordinate descriptions are listed in the Tables I–IV. Our calculations indicate that the molecular vibrations of NiTPC show some similarity to nickel meso-tetraphenylporphyrin (NiTPP), the porphyrin counterpart of NiTPC. The frequencies in the  $1440\text{--}1600\text{ cm}^{-1}$  region involve mainly the stretches of  $C_\beta C_\beta$ ,  $C_\alpha C_m$  bonds of corrole ring and CC bonds of phenyl groups. The  $C_\beta H$  in-plane bending and several pyrrole  $C_\alpha C_\beta / C_\alpha N$  stretches are located in the  $1040\text{--}1280\text{ cm}^{-1}$  region. The frequencies lower than  $970\text{ cm}^{-1}$  are mainly due to the out-of-plane vibrations of corrole skeleton and phenyl groups. On the other hand, IR and Raman spectra of NiTPC are more complex compared with NiTPP. This is expectable since the symmetry breaking of NiTPC causes splitting of the degenerated vibrational states and relaxation of IR and Raman selection rules. In addition, the vibrational coupling between the phenyl groups and the macrocycle is stronger in NiTPC than that in NiTPP, which results from the quite sharp dihedral angles between the  $C_m$ -phenyl groups and the corrole macrocycle ( $49.6^\circ$ ). In the following sections, assignments of the observed IR/Raman bands of NiTPC to local coordinates are discussed.

## B. Vibrational spectra of NiTPC

Figure 2 shows the experimental IR spectra of NiTPC in KBr pellets, in which the theoretical IR spectrum by B3LYP calculations is also provided (harmonic frequencies scaled with 0.97). It can be seen from Fig. 2 that the experimental and the theoretical spectra are coincident with each other in general. By comparing with the DFT-calculated band positions and intensities, the observed IR bands of NiTPC and other three complexes are assigned as listed in Table V. While both A and B modes are IR active, most of the strong IR absorptions of NiTPC in the high frequency region ( $900\text{--}1600\text{ cm}^{-1}$ ) belong to the A symmetry. Most of the IR bands with frequencies larger than  $900\text{ cm}^{-1}$  come from the in-plane corrole skeletal vibrations while the out-of-plane skeletal vibrations and phenyl modes dominate the region of  $400\text{--}900\text{ cm}^{-1}$ . In  $1400\text{--}1650\text{ cm}^{-1}$ , IR bands were measured at 1578, 1530, 1516, 1497, 1464, and  $1440\text{ cm}^{-1}$ . Theoretically, IR bands were predicted at 1608, 1582, 1523, 1496, 1465, and  $1435\text{ cm}^{-1}$ . The  $1578\text{ cm}^{-1}$  band can be as-

TABLE I Calculated in-plane skeletal modes (frequencies in  $\text{cm}^{-1}$ ) of NiTPC<sup>a</sup>.

A symmetry			B symmetry		
No.	Frequency	Mode description <sup>b</sup>	No.	Frequency	Mode description <sup>b</sup>
$\nu_1$	3191	$\nu(\text{C}_\beta^{\text{II}}\text{H})_s$	$\nu_{35}$	3191	$\nu(\text{C}_\beta^{\text{II}}\text{H})_s$
$\nu_2$	3180	$\nu(\text{C}_\beta^{\text{I}}\text{H})_s$	$\nu_{36}$	3180	$\nu(\text{C}_\beta^{\text{I}}\text{H})_s$
$\nu_3$	3175	$\nu(\text{C}_\beta^{\text{II}}\text{H})_{as}$	$\nu_{37}$	3175	$\nu(\text{C}_\beta^{\text{II}}\text{H})_{as}$
$\nu_4$	3160	$\nu(\text{C}_\beta^{\text{I}}\text{H})_{as}$	$\nu_{38}$	3159	$\nu(\text{C}_\beta^{\text{I}}\text{H})_{as}$
$\nu_5$	1532	$\nu(\text{C}_\alpha^{\text{I}}\text{C}_\alpha^{\text{I}})+\nu(\text{C}_\alpha^{\text{I,II}}\text{C}_m)_s$	$\nu_{39}$	1518	$\nu(\text{C}_\alpha^{\text{I,II}}\text{C}_m)_{as}$
$\nu_6$	1523	$\nu(\text{C}_\alpha^{\text{I,II}}\text{C}_m)_{as}+\nu(\text{Pyr-I}/2)_{as}$	$\nu_{40}$	1482	$\nu(\text{C}_\beta^{\text{I}}\text{C}_\beta^{\text{I}})+\nu(\text{C}_\beta^{\text{II}}\text{C}_\beta^{\text{II}})$
$\nu_7$	1499	$\nu(\text{C}_\alpha^{\text{II}}\text{C}'_m)_s$	$\nu_{41}$	1465	$\nu(\text{C}_\beta^{\text{II}}\text{C}_\beta^{\text{II}})+\nu(\text{C}_\beta^{\text{I}}\text{C}_\beta^{\text{I}})$
$\nu_8$	1459	$\nu(\text{C}_\beta^{\text{I}}\text{C}_\beta^{\text{I}})$	$\nu_{42}$	1459	$\nu(\text{C}_\alpha^{\text{II}}\text{C}'_m)_{as}$
$\nu_9$	1435	$\nu(\text{C}_\beta^{\text{II}}\text{C}_\beta^{\text{II}})$	$\nu_{43}$	1428	$\nu(\text{C}_\alpha^{\text{I,II}}\text{C}_m)_s$
$\nu_{10}$	1416	$\nu(\text{C}_\alpha^{\text{I,II}}\text{C}_m)_s$	$\nu_{44}$	1371	$\nu(\text{Pyr-I}/4)$
$\nu_{11}$	1374	$\nu(\text{Pyr-II}/2)_s+\nu(\text{Pyr-I}/2)_s$	$\nu_{45}$	1337	$\nu(\text{Pyr-I}/2)_s+\nu(\text{Pyr-II}/4)$
$\nu_{12}$	1345	$\nu(\text{Pyr-I}/4)$	$\nu_{46}$	1320	$\nu(\text{Pyr-I}/4)+\nu(\text{Pyr-II}/2)_{as}$
$\nu_{13}$	1333	$\nu(\text{Pyr-I}/2)_s$	$\nu_{47}$	1316	$\nu(\text{Pyr-II}/2)_s$
$\nu_{14}$	1297	$\nu(\text{Pyr-II}/4)$	$\nu_{48}$	1262	$\nu(\text{C}_m\text{C}_1)+\nu(\text{Pyr-II}/4)$
$\nu_{15}$	1276	$\delta(\text{C}_\beta^{\text{I}}\text{H})_{as}$	$\nu_{49}$	1221	$\delta(\text{C}_\beta^{\text{I}}\text{H})_{as}$
$\nu_{16}$	1258	$\nu(\text{C}_m\text{C}_1)+\nu(\text{Pyr-II}/4)$	$\nu_{50}$	1154	$\nu(\text{Pyr-I}/2)_{as}$
$\nu_{17}$	1224	$\nu(\text{C}'_m\text{C}_1)$	$\nu_{51}$	1132	$\delta(\text{C}_\beta^{\text{I}}\text{H})_{as}$
$\nu_{18}$	1189	$\delta(\text{C}_\beta^{\text{II}}\text{H})_{as}$	$\nu_{52}$	1072	$\delta(\text{C}_\beta^{\text{II}}\text{H})_s$
$\nu_{19}$	1082	$\delta(\text{C}_\beta^{\text{II}}\text{H})_s$	$\nu_{53}$	1049	$\delta(\text{C}_\beta^{\text{II}}\text{H})_s$
$\nu_{20}$	1059	$\delta(\text{C}_\beta^{\text{I}}\text{H})_s$	$\nu_{54}$	1025	$\nu(\text{Pyr-II br.})$
$\nu_{21}$	1028	$\nu(\text{Pyr-I}/2)_{as}$	$\nu_{55}$	1013	$\nu(\text{Pyr-II}/2)_{as}$
$\nu_{22}$	1017	$\nu(\text{Pyr-II br.})$	$\nu_{56}$	983	$\nu(\text{Pyr-I br.})$
$\nu_{23}$	1017	$\nu(\text{Pyr-II}/2)_{as}$	$\nu_{57}$	889	$\delta(\text{Pyr-I def.})_s$
$\nu_{24}$	984	$\nu(\text{Pyr-I br.})$	$\nu_{58}$	863	$\delta(\text{Pyr-I def.})_{as}+\delta(\text{Pyr-II def.})_s$
$\nu_{25}$	878	$\nu(\text{Pyr-II def.})_s+\nu(\text{Pyr-I def.})_{as}$	$\nu_{59}$	835	$\delta(\text{Pyr-I def.})_s+\delta(\text{Pyr-II def.})_{as}$
$\nu_{26}$	870	$\nu(\text{Pyr-II def.})_{as}$	$\nu_{60}$	806	$(\text{Pyr-I def.})_{as}+\delta(\text{Pyr-II def.})_{as}$
$\nu_{27}$	835	$\nu(\text{Pyr-I def.})_{as}+\nu(\text{Pyr-II def.})_{as}$	$\nu_{61}$	553	$\delta(\text{Pyr-I rot.})$
$\nu_{28}$	525	$\delta(\text{Pyr-I transl.})+\delta(\text{Pyr-II transl.})$	$\nu_{62}$	471	$\delta(\text{Pyr-II rot.})$
$\nu_{29}$	476	$\nu(\text{NiN})$	$\nu_{63}$	431	$\delta(\text{Pyr-I transl.})+\delta(\text{Pyr-II transl.})$
$\nu_{30}$	421	$\delta(\text{Pyr-I rot.})+\delta(\text{Pyr-II rot.})$	$\nu_{64}$	332	$\delta(\text{Pyr-II transl.})+\delta(\text{Pyr-I transl.})$
$\nu_{31}$	408	$\delta(\text{Pyr-II transl.})+\delta(\text{Pyr-I transl.})$	$\nu_{65}$	288	$\nu(\text{NiN})$
$\nu_{32}$	385	$\nu(\text{NiN})$	$\nu_{66}$	270	$\delta(\text{C}_m\text{-Ph})$
$\nu_{33}$	313	$\delta(\text{Pyr-I rot.})+\delta(\text{Pyr-II rot.})$	$\nu_{67}$	244	$\delta(\text{C}_m\text{-Ph})$
$\nu_{34}$	239	$\delta(\text{C}_m\text{-Ph})$			

<sup>a</sup> Calculated with B3LYP, basis sets: 6-31G\* for C, N, and H, TZVP for Ni. Frequency scaling factor is 0.97.

<sup>b</sup> For more than one local coordinates listed, the first one is dominate. “I” and “II” represent pyrrole ring I and II, respectively. “s”-symmetric, “as”-asymmetric, the description of local coordinates is according to Refs.[36, 37].

signed to the asymmetric CC stretching mode of phenyl substituents ( $\psi'_3$ ), which is calculated at  $1582\text{ cm}^{-1}$ . Another phenyl mode calculated at  $1608\text{ cm}^{-1}$  ( $\varphi_4$ ) is not observed in the experimental spectrum.

The stretching vibrations of  $\text{C}_\beta\text{C}_\beta$ ,  $\text{C}_\alpha\text{C}_m$ , and  $\text{C}_\alpha^{\text{I}}\text{C}_\alpha^{\text{I}}$  bonds are expected to appear in the  $1440\text{--}1650\text{ cm}^{-1}$  region due to their relatively large force constants. The  $1497\text{ cm}^{-1}$  band in the experimental spectrum is thought to correspond to the  $1496\text{ cm}^{-1}$  band in the theoretic spectrum and is assigned to  $\varphi_5$ , the CC stretch

and CH bending of the phenyl groups. The absorption band observed at  $1464\text{ cm}^{-1}$  is considered to correspond to the calculated band at  $1465\text{ cm}^{-1}$ , which according to the DFT calculation is a corrole skeletal mode ( $\nu_{41}$ ) involving the stretches of  $\text{C}_\beta^{\text{I}}\text{C}_\beta^{\text{I}}$  and  $\text{C}_\beta^{\text{II}}\text{C}_\beta^{\text{II}}$  bonds, and weakly, the  $\text{C}_\alpha^{\text{I,II}}\text{C}_m$  asymmetric stretch. The IR band observed at  $1440\text{ cm}^{-1}$  is thought to correspond to the calculated band at  $1435\text{ cm}^{-1}$  ( $\nu_9$ ). Calculations reveal that this band is due to the  $\text{C}_\beta^{\text{II}}\text{C}_\beta^{\text{II}}$  stretch that is

TABLE II Calculated out-of-plane skeletal modes (frequencies in  $\text{cm}^{-1}$ ) for NiTPC<sup>a</sup>.

A symmetry			B symmetry		
No.	Frequency	Mode description <sup>b</sup>	No.	Frequency	Mode description <sup>b</sup>
$\gamma_1$	881	$\gamma(\text{C}_{\beta}^{\text{II}}\text{H})_{\text{as}}$	$\gamma_{16}$	881	$\gamma(\text{C}_{\beta}^{\text{II}}\text{H})_{\text{as}}$
$\gamma_2$	854	$\gamma(\text{C}_{\beta}^{\text{I}}\text{H})_{\text{as}}$	$\gamma_{17}$	854	$\gamma(\text{C}_{\beta}^{\text{I}}\text{H})_{\text{as}}$
$\gamma_3$	773	$\gamma(\text{Pyr-II fold})_{\text{s}}$	$\gamma_{18}$	777	$\gamma(\text{Pyr-II fold})$
$\gamma_4$	762	$\gamma(\text{Pyr-I fold})_{\text{s}}$	$\gamma_{19}$	755	$\gamma(\text{Pyr-I fold})_{\text{s}}$
$\gamma_5$	706	$\gamma(\text{C}_{\beta}^{\text{II}}\text{H})_{\text{s}}$	$\gamma_{20}$	708	$\gamma(\text{C}_{\beta}^{\text{II}}\text{H})_{\text{s}}$
$\gamma_6$	700	$\gamma(\text{C}_{\beta}^{\text{I}}\text{H})_{\text{s}}$	$\gamma_{21}$	706	$\gamma(\text{C}_{\beta}^{\text{I}}\text{H})_{\text{s}}$
$\gamma_7$	656	$\gamma(\text{Pyr-II fold})_{\text{as}}+\gamma(\text{Pyr-I fold})_{\text{as}}$	$\gamma_{22}$	648	$\gamma(\text{Pyr-II fold})_{\text{as}}+\gamma(\text{Pyr-I fold})_{\text{as}}$
$\gamma_8$	647	$\gamma(\text{Pyr-I fold})_{\text{as}}+\gamma(\text{Pyr-II fold})_{\text{as}}$	$\gamma_{23}$	626	$\gamma(\text{Pyr-I fold})_{\text{as}}+\gamma(\text{Pyr-II fold})_{\text{as}}$
$\gamma_9$	562	$\gamma(\text{C}_{\alpha}^{\text{I,II}}\text{C}_{\text{m}})+\gamma(\text{C}_{\alpha}^{\text{II}}\text{C}'_{\text{m}})$	$\gamma_{24}$	569	$\gamma(\text{C}_{\alpha}^{\text{I,II}}\text{C}_{\text{m}})+\gamma(\text{C}_{\alpha}^{\text{II}}\text{C}'_{\text{m}})$
$\gamma_{10}$	365	$\gamma(\text{Pyr-I swivel})$	$\gamma_{25}$	564	$\gamma(\text{C}_{\alpha}^{\text{II}}\text{C}'_{\text{m}})$
$\gamma_{11}$	300	$\gamma(\text{Pyr-II swivel})$	$\gamma_{26}$	317	$\gamma(\text{Pyr-II swivel})$
$\gamma_{12}$	250	$\gamma(\text{Pyr tilt})$	$\gamma_{27}$	297	$\gamma(\text{Ni-N})$
$\gamma_{13}$	154	$\gamma(\text{Pyr tilt})$	$\gamma_{28}$	258	$\gamma(\text{Pyr-I swivel})$
$\gamma_{14}$	150	$\gamma(\text{C}_{\text{m}}\text{-Ph})$	$\gamma_{29}$	203	$\gamma(\text{Pyr tilt})$
$\gamma_{15}$	89	$\gamma(\text{Pyr transl.})$	$\gamma_{30}$	156	$\gamma(\text{Pyr tilt})$
			$\gamma_{31}$	113	$\gamma(\text{C}_{\text{m}}\text{-Ph})$
			$\gamma_{32}$	92	$\gamma(\text{C}_{\text{m}}\text{-Ph})$

<sup>a</sup> Calculated with B3LYP, basis sets: 6-31G\* for C, N, and H, TZVP for Ni. Frequency scaling factor is 0.97.

<sup>b</sup> For more than one local coordinates listed, the first one is dominate. “I” and “II” represent pyrrole ring I and II, respectively. “s”-symmetric, “as”-asymmetric, Refs.[36, 37] for the description of local coordinates.

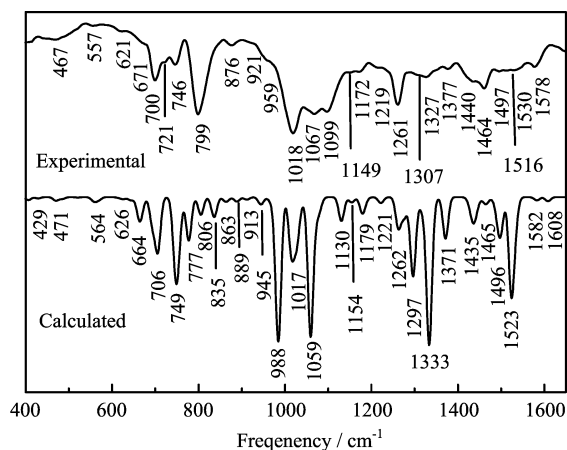


FIG. 2 Experimental (KBr disc) and calculated (frequency scaling factor is 0.97) IR absorption spectra of NiTPC.

strongly coupled to the  $\text{C}_{\beta}^{\text{I}}\text{C}_{\beta}^{\text{I}}$  stretch and symmetric  $\text{C}_{\alpha}^{\text{II}}\text{C}'_{\text{m}}$  stretch. Two weak bands were observed at 1530 and  $1516\text{ cm}^{-1}$  respectively. Theoretical calculation reveals a strong band at  $1523\text{ cm}^{-1}$  ( $\nu_6$ ). The  $\nu_6$  mode can be assigned to the asymmetric  $\text{C}_{\alpha}^{\text{I,II}}\text{C}_{\text{m}}$  stretch (A symmetry). The DFT calculation also manifests that this mode is strongly mixed with the half-ring asymmetric stretch of Pyr-I.

In the region of  $900\text{--}1440\text{ cm}^{-1}$ , the strongest ab-

sorption was observed at  $1018\text{ cm}^{-1}$  in the IR spectrum of NiTPC. However, DFT calculation predicts two strong IR bands at 1017 and  $988\text{ cm}^{-1}$ , respectively. The broad  $1017\text{ cm}^{-1}$  band contains IR absorptions of four skeletal vibrations, *i.e.*,  $\nu_{21}$ ,  $\nu_{22}$ ,  $\nu_{23}$ , and  $\nu_{55}$ , with  $\nu_{22}$  dominating IR intensity. The calculated  $988\text{ cm}^{-1}$  can be assigned to  $\varphi_8$ , the breathing vibration of the  $\text{C}_{\text{m}}$ -phenyl groups. We find that in this mode the vibrational coupling between the  $\text{C}_{\text{m}}$ -phenyl groups and the corrole skeleton is very severe. Since several strong IR absorptions are predicted in nearby positions, it is difficult to give a definite assignment of the observed  $1018\text{ cm}^{-1}$  band. We tend to attributing it to either  $\nu_{22}$  or  $\varphi_8$ .

Other strong absorption bands were observed at 1261, 1099, and  $1067\text{ cm}^{-1}$ . The  $1262\text{ cm}^{-1}$  band is thought to correspond to the calculated band at  $1262\text{ cm}^{-1}$  ( $\nu_{48}$ ), which is assigned to  $\text{C}_{\text{m}}\text{-C}_1$  stretch that is weakly coupled quarter-ring stretch of Pyr-II. The  $1099\text{ cm}^{-1}$  band is assigned to phenyl in-plane CH bending  $\psi_8$  that is calculated at  $1080\text{ cm}^{-1}$  as a shoulder band. The  $1067\text{ cm}^{-1}$  band in the experimental spectrum is thought due to  $\nu_{20}$  (symmetric bending of  $\text{C}_{\beta}^{\text{I}}\text{H}$  bond) which is calculated at  $1059\text{ cm}^{-1}$  with a large IR intensity. Additional weak absorption bands were observed at 1377, 1327, 1307, 1219, 1172, 1149, 959, and  $921\text{ cm}^{-1}$ . The calculated frequencies for these bands are at 1371 ( $\nu_{44}$ ), 1333 ( $\nu_{13}$ ), 1297 ( $\nu_{14}$ ), 1221 ( $\nu_{49}$ ),

TABLE III Calculated in-plane phenyl modes for NiTPC<sup>a</sup>.

A symmetry			B symmetry		
No.	$f^b$	Md <sup>c</sup>	No.	$f^b$	Md <sup>c</sup>
$\varphi_1$	3121	$\nu(\text{CH})$	$\varphi'_1$	3121	$\nu(\text{CH})$
$\varphi''_1$	3120	$\nu(\text{CH})$	$\psi_1$	3116	$\nu(\text{CH})$
$\psi'_1$	3116	$\nu(\text{CH})$	$\psi''_1$	3115	$\nu(\text{CH})$
$\varphi_2$	3109	$\nu(\text{CH})$	$\varphi'_2$	3109	$\nu(\text{CH})$
$\varphi''_2$	3108	$\nu(\text{CH})$	$\psi'_2$	3098	$\nu(\text{CH})$
$\psi_2$	3098	$\nu(\text{CH})$	$\psi''_2$	3098	$\nu(\text{CH})$
$\varphi_3$	3090	$\nu(\text{CH})$	$\varphi'_3$	3090	$\nu(\text{CH})$
$\varphi''_3$	3090	$\nu(\text{CH})$	$\varphi'_4$	1607	$\nu(\text{CC})$
$\varphi_4$	1608	$\nu(\text{CC})$	$\psi_3$	1583	$\nu(\text{CC})$
$\varphi'_4$	1607	$\nu(\text{CC})$	$\psi''_3$	1582	$\nu(\text{CC})$
$\psi'_3$	1582	$\nu(\text{CC})$	$\varphi_5$	1496	$\nu(\text{CC})+\delta(\text{CH})$
$\varphi'_5$	1495	$\nu(\text{CC})+\delta(\text{CH})$	$\psi_4$	1443	$\nu(\text{CC})+\delta(\text{CH})$
$\varphi''_5$	1492	$\nu(\text{CC})+\delta(\text{CH})$	$\psi''_4$	1440	$\nu(\text{CC})+\delta(\text{CH})$
$\psi'_4$	1440	$\nu(\text{CC})+\delta(\text{CH})$	$\psi_5$	1325	$\delta(\text{CH})$
$\psi''_5$	1325	$\delta(\text{CH})$	$\psi'_5$	1325	$\delta(\text{CH})$
$\psi_6$	1293	$\delta(\text{CH})$	$\psi'_6$	1292	$\delta(\text{CH})$
$\varphi_6$	1179	$\delta(\text{CH})$	$\psi''_6$	1287	$\delta(\text{CH})$
$\varphi''_6$	1178	$\delta(\text{CH})$	$\varphi_6$	1179	$\delta(\text{CH})$
$\psi'_7$	1158	$\delta(\text{CH})$	$\psi_7$	1158	$\delta(\text{CH})$
$\psi'_8$	1079	$\delta(\text{CH})$	$\psi''_7$	1157	$\delta(\text{CH})$
$\varphi_7$	1047	$\delta(\text{CH})$	$\psi_8$	1080	$\delta(\text{CH})$
$\varphi'_7$	1038	$\delta(\text{CH})$	$\psi''_8$	1078	$\delta(\text{CH})$
$\varphi'_8$	986	Ph. br.	$\varphi'_7$	1030	$\delta(\text{CH})$
$\varphi''_8$	986	Ph. br.	$\varphi_8$	988	Ph. br.
$\varphi_9$	670	$\delta(\text{CCC})$	$\varphi'_9$	664	$\delta(\text{CCC})$
$\varphi''_9$	643	$\delta(\text{CCC})$	$\psi_9$	616	$\delta(\text{CCC})$
$\psi'_9$	615	$\delta(\text{CCC})$	$\psi''_9$	615	$\delta(\text{CCC})$
$\varphi_{10}$	206	$\delta(\text{CCC})$	$\varphi'_{10}$	223	$\delta(\text{CCC})$
$\varphi''_{10}$	201	$\delta(\text{CCC})$	$\psi'_{10}$	32	Ph. rock.
$\psi_{10}$	51	Ph. rock.	$\psi''_{10}$	19	Ph. rock.

<sup>a</sup> Calculated with B3LYP, basis sets: 6-31G\* for C, N, and H, TZVP for Ni. Frequency scaling factor is 0.97.

<sup>b</sup>  $f$  denotes frequencies in  $\text{cm}^{-1}$ .

<sup>c</sup> Md denotes mode description.

1179 ( $\varphi'_6$ ), 1154 ( $\nu_{50}$ ), 945 ( $\sigma'_1$ ), and 913  $\text{cm}^{-1}$  ( $\pi_2$ ). According to DFT calculated atomic Cartesian displacements,  $\nu_{44}$ ,  $\nu_{13}$ ,  $\nu_{14}$ ,  $\nu_{49}$ , and  $\nu_{50}$  are the corrole skeletal modes, corresponding to the quarter-ring stretch of Pyr-I, symmetric half-ring stretch of Pyr-I, quarter-ring stretch of Pyr-II, asymmetric  $\text{C}'_{\beta}\text{H}$  in-plane bending, and asymmetric half-ring stretch of Pyr-I, respectively. The  $\varphi'_6$ ,  $\sigma'_1$ , and  $\pi_2$  are the phenyl modes which involve mainly the in-plane CH bending ( $\varphi'_6$ ) and out-of-plane CH wag ( $\sigma'_1$  and  $\pi_2$ ) of the phenyl groups.

In the low frequency region ( $<900 \text{ cm}^{-1}$ ), strong IR bands were measured at 799, 746, 721, and 700  $\text{cm}^{-1}$  while strong or middle strong absorptions were calculated at 835 ( $\nu_{59}$ ), 806 ( $\nu_{60}$ ), 777 ( $\gamma_{18}$ ), 749 ( $\pi_3$ ), and

TABLE IV Calculated out-of-plane phenyl modes for NiTPC<sup>a</sup>.

A symmetry			B symmetry		
No.	$f^b$	Md <sup>c</sup>	No.	$f^b$	Md <sup>c</sup>
$\pi_1$	968	$\gamma(\text{CH})$	$\pi'_1$	968	$\gamma(\text{CH})$
$\sigma_1$	945	$\gamma(\text{CH})$	$\pi''_1$	967	$\gamma(\text{CH})$
$\sigma''_1$	943	$\gamma(\text{CH})$	$\sigma'_1$	945	$\gamma(\text{CH})$
$\pi'_2$	913	$\gamma(\text{CH})$	$\pi_2$	913	$\gamma(\text{CH})$
$\sigma_2$	839	$\gamma(\text{CH})$	$\pi''_2$	911	$\gamma(\text{CH})$
$\sigma''_2$	837	$\gamma(\text{CH})$	$\sigma'_2$	839	$\gamma(\text{CH})$
$\pi_3$	749	$\gamma(\text{CH})$	$\pi'_3$	747	$\gamma(\text{CH})$
$\pi'_4$	694	$\gamma(\text{CCC})$	$\pi''_3$	743	$\gamma(\text{CH})$
$\pi''_5$	470	$\gamma(\text{CCC})$	$\pi_4$	693	$\gamma(\text{CCC})$
$\sigma'_3$	406	$\gamma(\text{CCC})$	$\pi'_4$	692	$\gamma(\text{CCC})$
$\sigma''_3$	404	$\gamma(\text{CCC})$	$\pi_5$	517	$\gamma(\text{CCC})$
$\sigma'_4$	56	Ph. spin.	$\pi'_5$	482	$\gamma(\text{CCC})$
$\pi'_6$	54	Ph. rock.	$\sigma_3$	407	$\gamma(\text{CCC})$
$\sigma''_4$	37	Ph. spin.	$\pi_6$	81	Ph. rock.
			$\sigma_4$	62	Ph. spin.
			$\pi''_6$	50	Ph. rock.

<sup>a</sup> Calculated with B3LYP, basis sets: 6-31G\* for C, N, and H, TZVP for Ni. Frequency scaling factor is 0.97.

<sup>b</sup>  $f$  denotes frequencies in  $\text{cm}^{-1}$ .

<sup>c</sup> Md denotes mode description.

706  $\text{cm}^{-1}$ . We tentatively assign the 799  $\text{cm}^{-1}$  band to  $\nu_{60}$ , the asymmetric deformation of Pyr-I which couples with asymmetric deformation of Pyr-II. It seems that  $\gamma_{18}$  may also contribute to the 799  $\text{cm}^{-1}$  band as it is a rather broad absorption. The  $\pi_3$  mode is due to the CH out-of-plane wag of the phenyl groups and is believed corresponding to the observed 746  $\text{cm}^{-1}$  band. The calculated strong absorption at 706  $\text{cm}^{-1}$  comes from the overlap of several normal modes, *i.e.*,  $\gamma_5$ ,  $\gamma_6$ ,  $\gamma_{20}$ , and  $\gamma_{21}$ , which belong to the out-of-plane wag of  $\text{C}'_{\beta}\text{H}$  and  $\text{C}''_{\beta}\text{H}$  bonds and are calculated at 708, 706, 706, and 700  $\text{cm}^{-1}$  respectively. The observed IR band at 700  $\text{cm}^{-1}$  is assigned to  $\gamma_{21}$  since this mode was predicted by calculations to have strongest intensity among the four modes. The observed shoulder at 721  $\text{cm}^{-1}$  is assigned to  $\gamma_{20}$  that is of a higher frequency than  $\gamma_{21}$ .

Weak IR bands were observed at 876, 671, 621, 557, and 467  $\text{cm}^{-1}$ , respectively. They are assigned to  $\nu_{58}$  (the deformation of pyrrole rings),  $\varphi'_9$ ,  $\varphi''_9$  (in plane deformation of phenyl groups),  $\gamma_{25}$  (out-of-plane wag of the  $\text{C}'_{\text{m}}$  atom), and  $\nu_{62}$  (the in-plane rotation of Pyr-II rings), respectively. These modes were calculated at 863, 664, 643, 564, and 471  $\text{cm}^{-1}$ , respectively.

Figure 3 compares the 488.0 nm excited Raman spectrum of NiTPC (solid powder) with the DFT calculated Raman spectrum. Observed Raman frequencies of NiTPC and other three complexes are listed in Table VI together with their assignments. It can be seen from it that most of the strong Raman bands belong to

TABLE V Observed IR frequencies (in  $\text{cm}^{-1}$ ) of MTPCs and their assignments. IR spectra were obtained with KBr discs.

	NiTPC	CuTPC	CoTPC	MnTPC		NiTPC	CuTPC	CoTPC	MnTPC
$\varphi'_4$		1593		1591	$\nu_{22}/\varphi_8$	1018	1017	1015	1022
$\psi'_3$	1578	1553	1578	1571	$\nu_{24}$		979	991	982
$\nu_{39}$	1530	1525	1541	1542	$\sigma'_1$	959		951	
$\nu_6$	1516	1515		1513	$\pi_2$	921	929	921	928
$\varphi_5$	1497	1490	1491	1490	$\nu_{57}$		890	889	
$\nu_{41}$	1464	1464	1464	1459	$\nu_{58}$	876	873	880	876
$\nu_9$	1440	1443	1444	1436	$\gamma_2$		842		841
$\nu_{43}$		1420	1418	1420	$\nu_{59}$		823	824	814
$\nu_{44}$	1377	1361	1371	1358	$\nu_{60}$	799	789	793	792
$\nu_{13}$	1327	1336	1350	1336	$\pi_3$	746	752	748	752
$\nu_{14}$	1307	1306	1312	1301	$\gamma_{20}$	721	714	718	717
$\nu_{48}$	1261	1267	1263	1282	$\gamma_{21}$	700	703	696	698
$\nu_{16}$		1239	1239	1260	$\varphi'_9$	671	672		668
$\nu_{49}$	1219		1214	1215	$\gamma_{22}$		654		650
$\varphi'_6$	1172	1174	1175	1177	$\varphi''_9$	621	617	619	623
$\nu_{50}$	1149	1155	1158	1152	$\gamma_{25}$	557	577	563	573
$\nu_{51}$		1120	1128		$\nu_{28}$		518	515	523
$\psi_8$	1099	1097	1099	1107	$\nu_{62}$	467	476	469	482
$\nu_{20}$	1067	1057	1059	1053	$\nu_{63}$		453		458

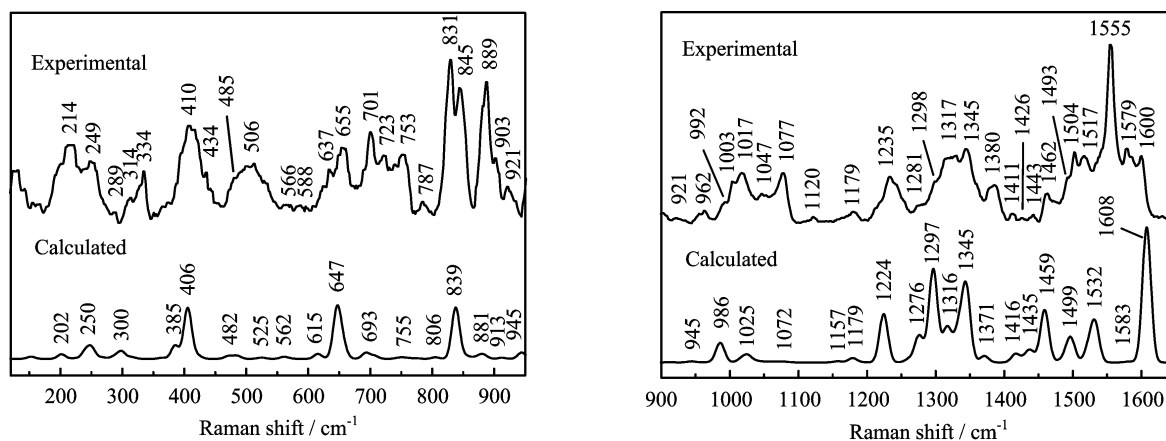


FIG. 3 Experimental (solid powder with 488 nm excitation) and calculated (frequency scaling factor is 0.97) Raman spectra of NiTPC.

A symmetry, although both A and B modes are Raman active.

As shown in Fig.3, the 488.0 nm excited spectrum of NiTPC gives rise to eight bands in 1440–1650  $\text{cm}^{-1}$  region, at 1600, 1579, 1555, 1517, 1504, 1493, 1462, and 1443  $\text{cm}^{-1}$ , respectively. The strongest 1555  $\text{cm}^{-1}$  band must be assigned to  $\nu_5$  ( $\text{C}_\alpha^{\text{I}}\text{C}_\alpha^{\text{I}}$  stretch coupled to the symmetric stretches of  $\text{C}_\alpha\text{C}_m$  bonds), since this mode is calculated to be the highest skeletal mode in this region and have a quite high Raman intensity. The 1517  $\text{cm}^{-1}$  band is assigned to  $\nu_6$  ( $\text{C}_\alpha^{\text{I,II}}\text{C}_m$  asymmetric stretch coupled to the asymmetric half-ring stretch of Pyr-I). The Raman band at 1462  $\text{cm}^{-1}$  is assigned to  $\nu_8$  ( $\text{C}_\beta^{\text{I}}\text{C}_\beta^{\text{I}}$  stretch with strong coupling to  $\text{C}_\alpha\text{C}_m$  sym-

metric stretch). We notice that the in-plane deformation of  $\text{C}_m$ -phenyls also contribute to this mode. The weak band at 1443  $\text{cm}^{-1}$  is better assigned to  $\nu_9$  ( $\text{C}_\beta^{\text{II}}\text{C}_\beta^{\text{II}}$  stretch in coupling with  $\text{C}_\alpha\text{C}_m$  symmetric stretch).

Several phenyl modes are observed in the Raman spectra with notable intensities. The observed strong 1600  $\text{cm}^{-1}$  band is assigned to the CC stretch of the phenyl groups ( $\varphi_4$ ), which was calculated to appear at 1608  $\text{cm}^{-1}$ . The middle-strong 1179  $\text{cm}^{-1}$  Raman band corresponds to the 1179  $\text{cm}^{-1}$  band in the theoretic spectrum and is assigned to the phenyl vibration involving mainly the phenyl CCH in-plane bending ( $\varphi_6$ ). The weak band at 1579  $\text{cm}^{-1}$  is thought to be a substituent mode involving the asymmetric CC stretch of

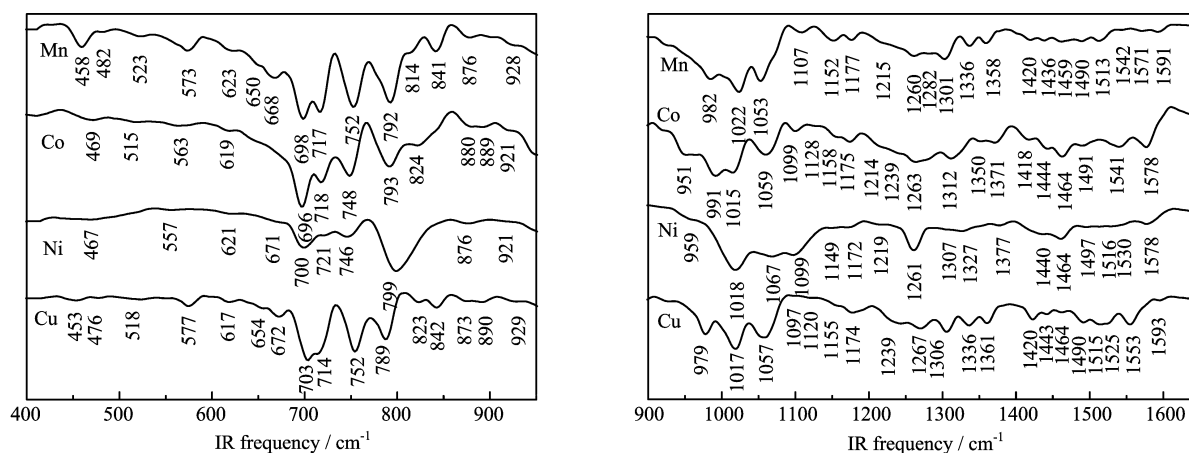


FIG. 4 Experimental IR absorption spectra (KBr discs) of CuTPC, NiTPC, CoTPC, and MnTPC.

phenyl groups ( $\psi_3$ ), but corrole skeletal  $C_\alpha C_m$  asymmetric stretch also contributes to this mode.

In the region of symmetric and asymmetric stretch of pyrrole rings ( $950\text{--}1440\text{ cm}^{-1}$ ), the strongest band was observed at  $1345\text{ cm}^{-1}$ . The corresponding band is calculated at the same position and is assigned to  $\nu_{12}$ , the quarter-ring stretch of Pyr-I conjoint with  $C_\alpha^{I,II}C_m$  symmetric stretch. The strong band at  $1235\text{ cm}^{-1}$  is thought to correspond to the calculated strong band at  $1224\text{ cm}^{-1}$ . This band ( $\nu_{17}$ ) is assigned to the  $C'_m C'_1$  stretch that is strongly coupled to the in-plane CCH bending of the phenyl groups. The observed  $1077\text{ cm}^{-1}$  band is assigned to the  $C_\beta^{II}H$  in-plane symmetric bending calculated at  $1072\text{ cm}^{-1}$ . The observed  $1380$ ,  $1317$ , and  $1017\text{ cm}^{-1}$  bands in this region are assigned to  $\nu_{44}$  (the quarter-ring stretch of Pyr-I),  $\nu_{47}$  (symmetric half-ring stretch of Pyr-II), and  $\nu_{54}$  (breathing of Pyr-II). The calculated frequencies of these modes are  $1371$  ( $\nu_{44}$ ),  $1316$  ( $\nu_{47}$ ), and  $1025\text{ cm}^{-1}$  ( $\nu_{54}$ ), respectively.

Additional weak or shoulder bands in this region were observed at  $1426$ ,  $1281$ ,  $1120$ ,  $1047$ ,  $1003$ , and  $992\text{ cm}^{-1}$ . The  $1426\text{ cm}^{-1}$  band is assigned to  $\nu_{43}$  (symmetric  $C_\alpha^{I,II}C_m$  stretch that is weakly coupled to the  $C_\beta^{II}C_\beta^{II}$  stretch). The  $1281\text{ cm}^{-1}$  band in the experimental spectrum is assigned to  $\nu_{15}$  (asymmetric  $C_\beta^I H$  in-plane bending coupled to the quarter-ring stretch of Pyr-I). The  $\nu_{43}$  and  $\nu_{15}$  modes are calculated as  $1428$  and  $1276\text{ cm}^{-1}$  with weak intensities. The weak bands at  $1120$  and  $1047\text{ cm}^{-1}$  in the experimental spectrum are assigned to  $\nu_{51}$  (the asymmetric  $C_\beta^I H$  in-plane bending) and  $\varphi_7$  (phenyl CH bending), respectively. Both  $\nu_{51}$  and  $\varphi_7$  have very low Raman intensities according to the calculation.

The Raman bands associating the pyrrole in-plane deformation, pyrrole rotational, pyrrole translational, pyrrole out-of-plane folding, pyrrole out-of-plane tilting and swiveling are expected to appear in the region lower than  $950\text{ cm}^{-1}$ . Quite strong Raman bands were observed at  $889$ ,  $845$ , and  $831\text{ cm}^{-1}$ . The  $889\text{ cm}^{-1}$  band

can be assigned to  $\gamma_1$  calculated at  $881\text{ cm}^{-1}$ , which is a skeletal mode involving the asymmetric  $C_\beta^{II}H$  out-of-plane wag. The  $845$  and  $831\text{ cm}^{-1}$  bands are thought to come from the CH out-of-plane wag of phenyl groups ( $\sigma_2$ ) and the in-plane deformation of the pyrroles ( $\nu_{27}$ ), respectively. The calculated frequencies of  $\sigma_2$  and  $\nu_{27}$  are  $839$  and  $835\text{ cm}^{-1}$ , respectively, and both modes are of considerable Raman intensities. The weak bands observed at  $921$  and  $903\text{ cm}^{-1}$  are thought to correspond to the calculated bands at  $913\text{ cm}^{-1}$  ( $\pi'_2$ ) and  $889\text{ cm}^{-1}$  ( $\nu_{57}$ ), respectively. The former is due to the out-of-plane CH wag of  $C_m$ -phenyl groups, while the later is a corrole skeletal mode associating with the in-plane deformation of Pyr-I. The weak band at  $787\text{ cm}^{-1}$  is assigned to  $\gamma_4$  calculated at  $762\text{ cm}^{-1}$  (symmetric out-of-plane folding of Pyr-I).

In the region lower than  $800\text{ cm}^{-1}$ , Raman bands observed at  $753$ ,  $655$ ,  $289$ , and  $249\text{ cm}^{-1}$  are considered to be the out-of-plane skeletal vibrations, which were predicted at  $755\text{ cm}^{-1}$  ( $\gamma_{19}$ ),  $647\text{ cm}^{-1}$  ( $\gamma_8$ ),  $300\text{ cm}^{-1}$  ( $\gamma_{11}$ ), and  $250\text{ cm}^{-1}$  ( $\gamma_{12}$ ), respectively. The  $\gamma_{19}$  and  $\gamma_8$  modes are assigned to the symmetric and asymmetric oop folding of Pyr-I, respectively. The  $\gamma_{11}$  mode is assigned to the swiveling of Pyr-II, while  $\gamma_{12}$  is due to the pyrrole tilting motions (*i.e.*, saddling of corrole macrocycle). The observed  $314\text{ cm}^{-1}$  band is assigned to the in-plane skeletal mode  $\nu_{33}$  (calculated at  $313\text{ cm}^{-1}$ ).

Raman bands at  $701$ ,  $485$ , and  $410\text{ cm}^{-1}$  in the experimental spectrum are assigned to phenyl out-of-plane CCC deformation vibrations  $\pi'_4$ ,  $\pi'_5$ , and  $\sigma'_3$ . The calculated frequencies of these modes are  $694$  ( $\pi'_4$ ),  $486$  ( $\pi'_5$ ), and  $406\text{ cm}^{-1}$  ( $\sigma'_3$ ), respectively.

### C. Comparison NiTPC with CoTPC, CuTPC and MnTPC

Figure 4 shows the infrared absorption spectra of MTPCs (M=Cu, Ni, Co, and Mn). It can be seen that the IR spectra of the four complexes are quite similar to each other. Some significant dif-



TABLE VI Observed Raman frequencies (in  $\text{cm}^{-1}$ ) of MTPCs and their assignments<sup>a</sup>.

	NiTPC	CuTPC	CoTPC	MnTPC		NiTPC	CuTPC	CoTPC	MnTPC
$\varphi_4$	1600	1596	1600	1594	$\sigma_1$	962		953	
$\psi_3$	1579	1559	1582	1572	$\pi'_2$	921		921	918
$\nu_5$	1555	1521	1564	1537	$\nu_{57}$	903		901	898
$\nu_6$	1517	1509	1534	1514	$\gamma_1$	889	888	885	882
$\nu_7$	1504	1495	1507	1493	$\sigma_2$	845	844	845	842
$\nu_{40}$	1493	1485			$\nu_{27}$	831	822	831	829/815
$\nu_8$	1462	1462	1473	1474	$\gamma_4$	787	783	785	794
$\nu_9$	1443	1438		1435	$\gamma_{19}$	753	747	747	751
$\nu_{43}$	1426	1425		1423	$\pi'_3$	723	722	723	716
$\nu_{10}$	1411	1404	1419		$\pi'_4$	701	698	699	698
$\nu_{44}$	1380	1363	1370	1362	$\gamma_8$	655	653	661	652
$\nu_{12}$	1345	1340		1347	$\gamma_{23}$	637	631	642	643
$\nu_{47}$	1317	1302	1328	1330	$\psi'_9$		618		625
$\nu_{14}$	1298	1292		1302	$\gamma_{24}$	588	589	592	576
$\nu_{15}$	1281	1263	1288	1263	$\gamma_9$	566	547	560	557
$\nu_{17}$	1235	1235	1236	1235	$\nu_{28}$	506	512	517	523
$\nu_{49}$		1218		1219	$\pi'_5$	485	470	450	457
$\varphi_6$	1179	1186	1186	1176	$\nu_{63}$	434		427	423
$\psi'_7$		1172		1151	$\sigma'_3$	410	412	406	406
$\nu_{51}$	1120			1116	$\nu_{32}$		402	391	382
$\nu_{52}$	1077	1072	1072	1074	$\nu_{64}$	334	358	334	345
$\varphi_7$	1047	1044	1042	1055	$\nu_{33}$	314	324		
$\nu_{54}$	1017	1019	1019	1021	$\gamma_{11}$	289	307		305
$\nu_{23}$	1003	1000		1002	$\gamma_{28}$			266	267
$\varphi'_8$	992	979	992	986	$\gamma_{12}$	249	245	243	249
					$\varphi''_{10}$	214	210	199	215

<sup>a</sup> Raman spectra obtained from solid powder. Excitation wavelength is 488 nm.

ferences in the IR spectra of MTPCs are as follows. The CC stretch of  $\text{C}_m$ -phenyls ( $\varphi'_4$ ) were clearly observed in CuTPC ( $1593 \text{ cm}^{-1}$ ) and MnTPC ( $1591 \text{ cm}^{-1}$ ), but it was masked by nearby strong band ( $1578 \text{ cm}^{-1}$ ) in NiTPC and CoTPC. The IR bands at  $1420/1418/1420$ ,  $979/991/982$ , and  $823/824/814 \text{ cm}^{-1}$  for CuTPC/CoTPC/MnTPC were assigned to  $\nu_{43}$  (the symmetric  $\text{C}_\alpha^{\text{I,II}}\text{C}_m$  stretch),  $\nu_{24}$  (breathing of Pyr-I), and  $\nu_{59}$  (pyrrole in-plane deformation), respectively. These bands were not observed or extreme weak in NiTPC. The strong band of NiTPC at  $799 \text{ cm}^{-1}$  is assigned to the  $\nu_{60}$  mode. This mode was observed at  $789/793/792 \text{ cm}^{-1}$  in CuTPC/CoTPC/MnTPC, showing evident shift for different complexes. The middle strong band at  $842/841 \text{ cm}^{-1}$  in CuTPC/MnTPC were not observed in NiTPC and CoTPC. The band is assigned to  $\gamma_2$ , the out-of-plane  $\text{C}_\beta^{\text{II}}\text{H}$  wag.

The 488 nm excited Raman spectra of MTPCs ( $\text{M}=\text{Cu, Ni, Co, Mn}$ ) are displayed in Fig.5. As shown in Fig.5, the relative intensities of several Raman bands vary for different metal complexes, which is probably due to the different resonance enhancement pattern

since the main electronic absorption of MTPCs, which varies from 415 nm to 440 nm, is not far from the laser line. In spite of this, Raman spectra of the four complexes show similar general trends and thus correlation of the corresponding bands of MTPCs can be made out as listed in Table VI. Some differences between the Raman spectra of the four complexes are noticed. For instance, the  $\nu_{40}$  was observed in NiTPC and CuTPC but not in CoTPC and MnTPC while  $\nu_{49}$  and  $\psi'_7$  were observed in CuTPC/MnTPC but not in NiTPC and CoTPC. The quality of the Raman spectrum of CoTPC is poorer as compared with others due to strong photolysis of the sample. Consequently,  $\nu_9$ ,  $\nu_{43}$ ,  $\nu_{12}$ ,  $\nu_{14}$ , and  $\nu_{23}$  of CoTPC were not observed due to the weak intensities.

The molecular structure changes generated by different metal substitutions can be reflected in the Raman spectra. For metalloporphyrins (MPors) and metallotetraazaporphyrins (MTAPs), empirical correlations have been found between the Raman frequencies and the core-size Ct-N (*i.e.*, the distance between the center of the porphyrin/tetraazaporphyrin ring and the pyr-

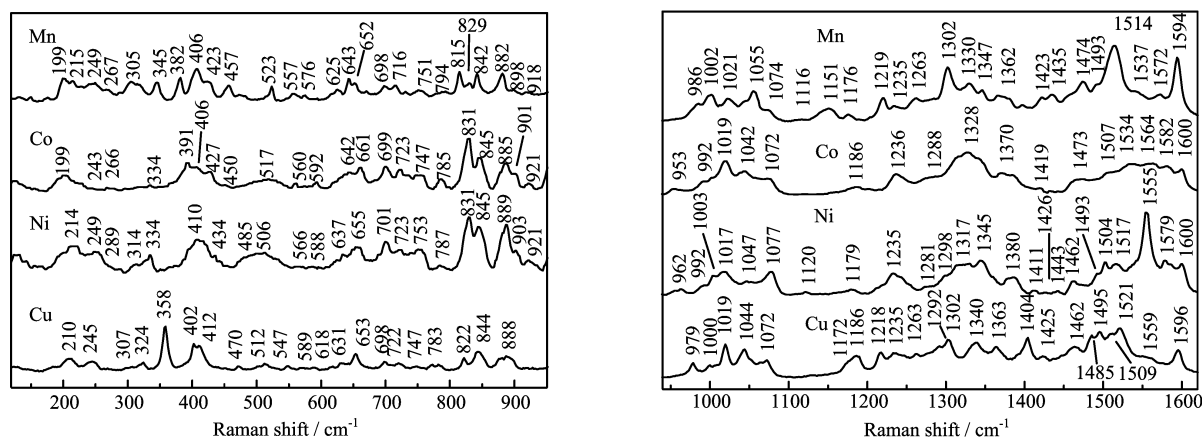


FIG. 5 Experimental Raman spectra (solid powder with 488 nm excitation) of CuTPC, NiTPC, CoTPC, and MnTPC.

role nitrogen atom) [25, 38, 39]. Especially, the Raman bands in the 1450–1600  $\text{cm}^{-1}$  region, which are usually due to the stretching vibrations of the  $\text{C}_\beta\text{C}_\beta$  and  $\text{C}_\alpha\text{C}_m$  bonds, have been found sensitive to the core-size change.

We have studied the frequency/core-size relationship of MTPC by correlating the some observed Raman frequencies in the 1450–1600  $\text{cm}^{-1}$  region with the DFT calculated M–N distances (averaged over four M–N bonds). Figure 6 displays the relationships between the core-sizes (M–N distances) and the frequencies of  $\nu_5$ ,  $\nu_6$ ,  $\nu_7$ , and  $\psi_3$  modes as well as the linear fitting of corresponding data. It is found that, in general, the Raman frequencies of these modes decrease with the increase of M–N distance. The slope of the fitted line reflects the relative sensitivity of a specific vibrational frequency to the core-size change. We found that with the decrease of core-size by 0.01 Å, the frequencies of  $\nu_5$ ,  $\nu_6$ ,  $\nu_7$ , and  $\psi_3$  increase by 6.76, 3.31, 2.23, and 3.48  $\text{cm}^{-1}$ , respectively.

Among these four modes,  $\nu_5$  has the highest sensitivity and best linearity for its frequency/core-size relationship and thus can be used as a structure mark band for the corrole compounds. This mode is assigned to  $\text{C}_\alpha^{\text{I}}\text{C}_\alpha^{\text{I}}$  stretch which is strongly coupled to the symmetric  $\text{C}_\alpha^{\text{I,II}}\text{C}_m$  and  $\text{C}_\alpha^{\text{II}}\text{C}'_m$  stretches. The  $\nu_6$  and  $\nu_7$  are assigned to the asymmetric  $\text{C}_\alpha^{\text{I,II}}\text{C}_m$  stretch and symmetric  $\text{C}_\alpha^{\text{II}}\text{C}'_m$  stretch, respectively. As a phenyl mode, the core-size dependent of  $\psi_3$  is unusual. We find that while  $\psi_3$  is nominally assigned to a phenyl mode, it contains a considerable component of asymmetric stretches of  $\text{C}_\alpha^{\text{I,II}}\text{C}_m$  and  $\text{C}_\alpha^{\text{II}}\text{C}'_m$ . We consider that this vibrational coupling between substituents and corrole skeleton is partly responsible for the core-size dependent of  $\psi_3$ . In view that the  $\text{C}_\alpha^{\text{I}}\text{C}_\alpha^{\text{I}}$  and  $\text{C}_\alpha\text{C}_m$  stretches are sensitive to the core-size change, it can be deduced that the corrole macrocycle accommodates the metal ions with different sizes by adapting the  $\text{C}_\alpha^{\text{I}}\text{C}_\alpha^{\text{I}}$  and  $\text{C}_\alpha\text{C}_m$  bonds.

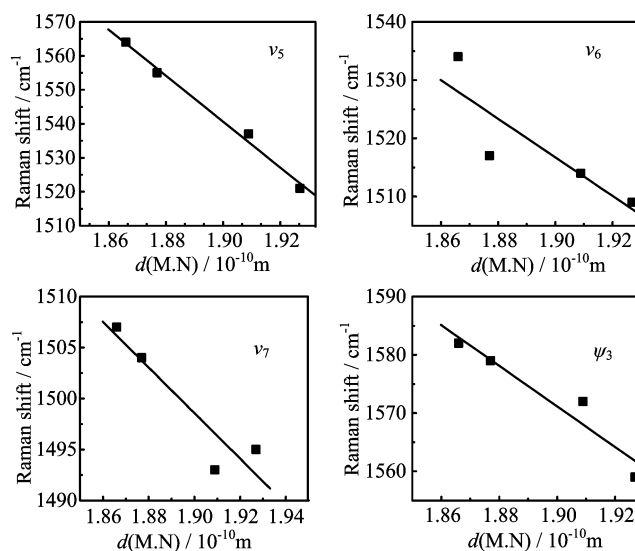


FIG. 6 Correlation between the Raman frequencies (for  $\nu_5$ ,  $\nu_6$ ,  $\nu_7$ , and  $\psi_3$  modes) and the calculated core-sizes of MTPC (M=Cu, Ni, Co, Mn). Frequencies from solid spectra excited at 488.0 nm.

#### IV. CONCLUSION

We have studied the IR absorption and 488 nm excited Raman scattering spectra of serious metal complexes of triphenylcorrole (MTPCs, M=Cu, Ni, Co, Mn). Density functional theory calculation using B3LYP functional were carried out to reveal the ground-state structures and vibrational states of MTPCs. Assignments of the observed IR and Raman bands of MTPCs to local coordinates were proposed based on the DFT calculations. Due to the symmetry lowering, the vibrational spectra of MTPCs are much more complex as compared with their porphyrin counterparts MTPP. The relationship between the Raman/IR frequencies and the structure of corrole ring was investigated. It was found that the vibrations involving the  $\text{C}_\alpha^{\text{I}}\text{C}_\alpha^{\text{I}}$  stretch and  $\text{C}_\alpha\text{C}_m$

stretch are sensitive to the size of corrole core. The frequency of  $\nu_5$  increases linearly with the decrease of the corrole core-sizes and may be used as a mark to evaluate the structural change of the metalloporphyrins.

## V. ACKNOWLEDGMENTS

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